

# The effect of pressure on the electrical resistance of lithium, sodium and potassium at low temperatures

BY J. S. DUGDALE

*Division of Pure Physics, National Research Council, Ottawa*

AND D. GUGAN

*H. H. Wills Physics Laboratory, University of Bristol*

(Communicated by D. K. C. MacDonald, F.R.S.—Received 2 April 1962)

Measurements have been made of the electrical resistivity of lithium, sodium and potassium at temperatures between 2 and 300 °K and at pressures up to 3000 atm. From our results we have calculated the ideal electrical resistivity,  $\rho_i$ , and its volume derivative as functions of temperature for conditions of constant density. It is shown that, as predicted by simple theory, there is a linear relation between the temperature and volume coefficients of  $\rho_i$  for each metal. We conclude that the magnitude of the volume coefficient of  $\rho_i$  does not, at high temperatures at least, agree with present theoretical predictions and that this coefficient is closely connected with the high-temperature value of the thermoelectric power.

## 1. INTRODUCTION

In order to obtain a general understanding of how the effect of pressure on the electrical resistivity of a pure metal changes with temperature it is convenient to make use of the following simple expression for the ideal electrical resistivity of a metal:

$$\rho_i = \frac{KT}{M\theta_R^2} f(T/\theta_R). \quad (1)$$

$\theta_R$  is here a constant, having the dimensions of temperature, which characterizes the resistive properties of the metal,  $M$  is the mass of the metallic ions, and  $K$  is a parameter which measures the interaction between the conduction electrons and the lattice vibrations.  $f$  is a function which becomes constant at high temperatures and which at very low temperatures is expected to vary as  $(T/\theta_R)^4$ . One example of such a function occurs in the Bloch-Grüneisen expression for the temperature dependence of the ideal resistivity of a metal, but for our present purposes we do not need to make any assumption about the form of  $f$  except that it is independent of volume. We emphasize, however, that  $K$  and  $\theta_R$  are assumed to be independent of temperature and to depend only on the volume.

Under these conditions, the *volume* coefficient of the ideal resistivity is related to the temperature coefficient of the ideal resistivity in the following way:

$$\left(\frac{\partial \ln \rho_i}{\partial \ln V}\right)_T = \frac{d \ln K}{d \ln V} - \frac{d \ln \theta_R}{d \ln V} \left\{ 1 + \left(\frac{\partial \ln \rho_i}{\partial \ln T}\right)_V \right\}. \quad (2)$$

At high temperatures ( $T \gtrsim \theta$ )  $\partial \ln \rho_i / \partial \ln T$  tends to unity for most metals (at least at constant density) so that in this region we may write:

$$\partial \ln \rho_i / \partial \ln V = (d \ln K / d \ln V) + 2\gamma_R, \quad (3)$$